

SFB-lectures

Phenomenological Aspects of String Theory

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Outline of the lectures:

1. Overview (January 19)
2. Compactifications (January 26, today)
3. Supersymmetry Breaking (February 2)

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String Theory

basic idea: point-like objects \rightarrow extended objects (strings)

Strings move in 10-dimensional space-time background

one way to make contact with “our world”: Compactification

\Rightarrow space-time background:

$$\mathcal{M}_{10} = R_{1,3} \times Y_6$$

$R_{1,3}$: four-dimensional Minkowski-space

Y_6 : compact manifold – determines amount of supersymmetry

alternative: Y_6 is not a manifold (but a conformal field theory).

Different string theories:

Type I, Type II, Heterotic

they differ in spectrum of excitations and their interactions

massless spectrum in $R_{1,9}$:

	Type IIA	Type IIB	Type I	Heterotic
B:	$G_{MN}, H_3 = dB_2, \Phi$			
B:	$F_2 = dC_1, F_4 = dC_3$	$l, F_3 = dC_2, F_5^* = dC_4$	$F_2^a = DA_1^a, G = SO(32), E_8 \times E_8$	
F	$\Psi^{1,2}, \lambda^{1,2}$	$\Psi^{1,2}, \lambda^{1,2}$	$\Psi^1, \lambda^1, \chi^a$	

$F_p = p$ -form field strength

$A, B, C =$ gauge potentials

$\Psi, \lambda, \chi =$ gravitino, dilatino, gaugino

Compactification: determine Y_6

Lorentz group on space-time background $\mathcal{M}_{10} = R_{1,3} \times Y_6$ decomposes

$$SO(1,9) \rightarrow SO(1,3) \times SO(6)$$

spinor decompose accordingly:

$$\mathbf{16} \rightarrow (\mathbf{2}, \mathbf{4}) \oplus (\bar{\mathbf{2}}, \bar{\mathbf{4}})$$

existence of supercharges Q requires:

\Rightarrow nowhere vanishing and globally defined spinor η needs to exist

$\Rightarrow Y_6$ has to be manifold with G -structure

\Rightarrow structure group is reduced $SO(6) \rightarrow G$

example:

$$SO(6) \rightarrow SU(3) : \quad \mathbf{4} \rightarrow \mathbf{3} + \mathbf{1}$$

$\Rightarrow Y_6$ has $SU(3)$ -structure

type II: $N = 2$ in $d = 4$, heterotic, type I: $N = 1$ in $d = 4$

supersymmetry of the background requires: $\langle \delta \text{fermions} \rangle = 0$

$$\delta \Psi_m = \nabla_m \eta + (\gamma F)_m \cdot \eta + \dots = 0, \quad m = 1, \dots, 6$$

$$\delta \lambda = (\gamma F) \cdot \eta + \dots = 0,$$

$$\delta \chi^a = (\gamma F^a) \cdot \eta + \dots = 0,$$

$\langle F \rangle \neq 0$ corresponds to background flux

\Rightarrow for $F = 0$: $\nabla_m \eta = 0 \quad \Rightarrow Y_6$ is Calabi-Yau manifold

\Rightarrow for $F \neq 0$: $\nabla_m \eta \neq 0 \quad \Rightarrow Y_6$ is manifold with $SU(3)$ structure

heterotic and type I: additional constraint

$$\text{tr}(R \wedge R) = (F^a \wedge F^a) \leftrightarrow R_{[mn]p}{}^q R_{rs]q}{}^p = F_{[mn}^a F_{rs]}^a \neq 0$$

$\Rightarrow F_{mn}^a$ needs a VEV

\Rightarrow breaking of the gauge group,

$$\begin{aligned} \text{e.g.: standard embedding} \quad E_8 \times E_8 &\rightarrow E_8 \times E_6 \\ &SO(32) \rightarrow SO(26) \end{aligned}$$

physical interpretation: E_6 is GUT group, E_8 is hidden sector

Kaluza-Klein compactification in space-time background: $R_{1,3} \times Y_6$

- massless scalars

$$\Delta_{10}\phi = (\Delta_4 + \Delta_6)\phi = (\Delta_4 + m^2)\phi = 0$$

for

$$\phi = \sum_n \phi^n(x) Y^n(y) \quad \text{and} \quad \Delta_6 Y^n = m_n^2 Y^n$$

- similarly for p-forms C_p and the metric $g_{mn} = g_{mn}^0 + \delta g_{mn}$
 \Rightarrow massless $d = 4$ spectrum = zero modes of Δ_6 = harmonic forms on Y_6
- harmonic forms $\omega_{(p,q)}$ are elements of cohomology groups $H^{(p,q)}(Y)$
multiplicity is counted by Hodge numbers
for Calabi-Yau manifold:

$$h^{p,q} = \dim H^{p,q}(Y),$$

			1		
		0		0	
	0		$h^{1,1}$		0
1		$h^{1,2}$		$h^{1,2}$	1
	0		$h^{1,1}$		0
		0		0	
			1		

massless spectrum arranges itself into supermultiplets

- heterotic string (with standard embedding and $G = E_8 \times E_6$)
 - gravitational multiplet
 - vector multiplets in adjoint of $E_8 \times E_6$
 - chiral matter multiplets

$$h^{(1,1)}(\overline{\mathbf{27}} + \mathbf{1}) \oplus h^{(1,2)}(\mathbf{27} + \mathbf{1})$$

\Rightarrow # of chiral generations $\sim h^{(1,1)} - h^{(1,2)} \sim \chi$ (Euler number)

- type IIA:

- gravitational multiplet
- $h^{(1,1)}$ $U(1)$ -vector multiplets
- $h^{(1,2)}$ uncharged hypermultiplets

\Rightarrow type II does not look very promising

\Rightarrow need to insert space-time filling D-branes

\Rightarrow non-Abelian gauge bosons and charged matter as excitations of D-branes

which choices went in so far?

- which string theory
- which compactifications manifolds Y_6
- which embedding
- which D-branes

$\Rightarrow \mathcal{O}(10^5)$ choices (far more if non-geometrical backgrounds included)

did not impose:

- $n_g = 3$
- constraint on G /spectrum

The low effective action

Integrating out the heavy modes H results in the low energy effective action $\mathcal{L}_{\text{eff}}(L)$ of the light modes L

The diagram shows a four-point contact interaction (a circle with four external lines labeled L) equal to the sum of two exchange processes:

- An s-channel exchange process where two L lines meet at a vertex, exchange a L particle, and then split into two L lines.
- A t-channel exchange process where two L lines meet at a vertex, exchange a H particle, and then split into two L lines.

+ t and u channels

For $p^2 \ll M_{\text{string}}^2$, the s-channel exchange process is approximated by:

- The s-channel exchange process with a L particle in the internal line.
- A contact interaction represented by a central black dot where four L lines meet.

+ t and u channels

for $p^2 \ll M_{\text{string}}^2$: amplitudes of an effective field theory

$D = 4, N = 1$ effective Lagrangian

[Wess, Bagger]

⇒ spectrum:

multiplet	B	F
gravity multiplet	$g_{\mu\nu}$	Ψ_μ
vector multiplets	V_μ	χ
chiral multiplets	M	λ

⇒ effective Lagrangian

$$L = -\left(\frac{1}{2}R + G_{I\bar{J}}D_\mu M^I D^\mu \bar{M}^{\bar{J}} + \frac{1}{4}\text{Re}f_{\kappa\lambda} F_{\mu\nu}^\kappa F^{\lambda\mu\nu} + \frac{1}{4}\text{Im}f_{\kappa\lambda} F^\kappa \tilde{F}^\lambda + V\right),$$

+ fermions

$$V = e^K (G^{I\bar{J}} D_I W D_{\bar{J}} \bar{W} - 3|W|^2) + \frac{1}{2} (\text{Re } f)^{-1 \kappa\lambda} D_\kappa D_\lambda .$$

L is completely determined in terms of K, W, f :

- **Kähler metric:** $G_{I\bar{J}} = \partial_I \bar{\partial}_{\bar{J}} K(M, \bar{M})$
- **holomorphic superpotential:** $W(M), \quad D_I W = \partial_I W + (\partial_I K)W$
- **holomorphic gauge kinetic function:** $f(M) = g^{-2} + i \frac{\Theta}{8\pi^2}$

Properties of K, W, f

- Kähler potential K receives quantum corrections at all orders
- Superpotential W has no perturbative quantum corrections

$$W = W^{(0)} + W^{(\text{np})}$$

- gauge kinetic function f has only 1-loop quantum corrections

$$f = f^{(0)} + f^{(1)} + f^{(\text{np})}$$

Next: compute K, W, f for your favorite string background.

Generic features perturbatively:

- ⇨ effective potential V has flat directions (moduli T)
 - ⇒ continuous vacuum degeneracy parameterized by $\langle T \rangle$
- ⇨ $W(T, Q) = Y_{ijk}(T) Q^i Q^j Q^k$, Q = charged matter field.
 - ⇒ Yukawa coupling Y dynamically determined!
- ⇨ similarly, gauge couplings $g = g(\langle T \rangle)$
 - ⇒ g is dynamically determined
 - If moduli are undetermined couplings are still free parameters.
- ⇨ need to stabilize moduli and break supersymmetry
 - known mechanism:
 - 1) Fluxes/generalized geometries
 - 2) Gaugino condensation/non-perturbative effects

Background fluxes [Polchinski, Strominger,...]

allow $\int_{\gamma_p^I \in Y} F_p = e_I \neq 0$ keeping $dF_p = 0 = d^\dagger F_p$

$$\Rightarrow F_p = e_I \omega_p^I, \quad \omega_p \in H^p(Y)$$

$$e_I = \text{const.} = \begin{cases} \text{quantized in string theory} \\ \text{continuous in low energy approximation} \end{cases}$$

consistency: tadpole cancellation condition

properties:

- large Y : e_I small perturbation such that light spectrum does not change
- low energy supergravity \Rightarrow gauged/massive supergravity
- potential generated \Rightarrow vacuum degeneracy (partially) lifted
- supersymmetry spontaneously broken

Background fluxes and the landscape

Fluxes are additional discrete parameters which are allowed in a compactification.

⇒ How many fluxes do we have?

number of cycles: $I = 1, \dots, \mathcal{O}(100)$

allowed values of e_I (solutions of tadpole condition): $e = 1, \dots, 10$.

⇒ for a given Calabi-Yau we have $\approx 10^{300}$ different vacua (landscape of vacua)

⇒ what do they determine?

generate potential for moduli ⇒ $\langle T \rangle$

- value of the moduli ⇒ value of Yukawa and gauge couplings
- spontaneous supersymmetry breaking
- value of the cosmological constant [Bousso, Polchinski]

$$\Lambda = -\Lambda_0 + \sum_{IJ} e_I e_J N^{IJ}$$

⇒ anthropic/statistical explanation of small Λ :

$$\exists e_I : 0 < \Lambda < \epsilon$$

Discussion/Summary

- ⇨ can construct realistic string backgrounds in $d = 4$ via compactifications
- ⇨ can compute the effective action reliably in perturbation theory
- ⇨ background fluxes appear as additional parameters
 - ⇒ landscape of vacua
- ⇨ open question: are they all consistent non-pertubatively?

(are they all vacua of M-theory?)

- if no, which are consistent?
- if yes, what does it mean?

one option:

